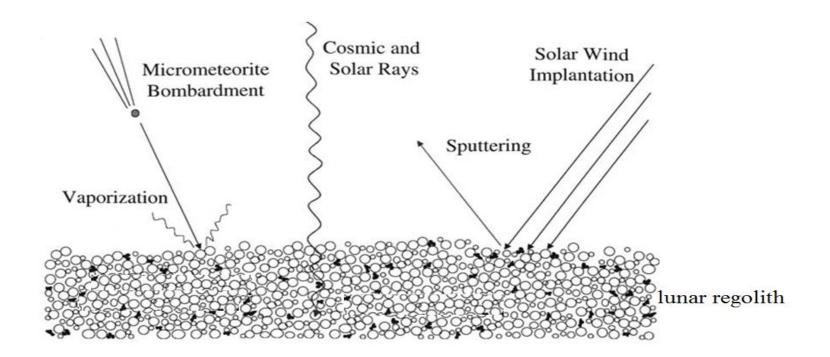
# Sputtering of Lunar Regolith by Solar Wind Protons and Heavy Ions

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### Introduction

- Lunar surface material is accessible to the space weathering factors
- Solar wind protons and heavy ions with kinetic energies of about 1 keV/amu

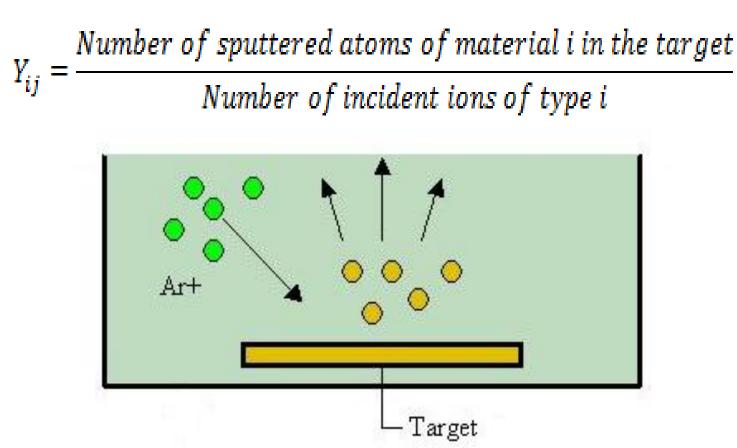


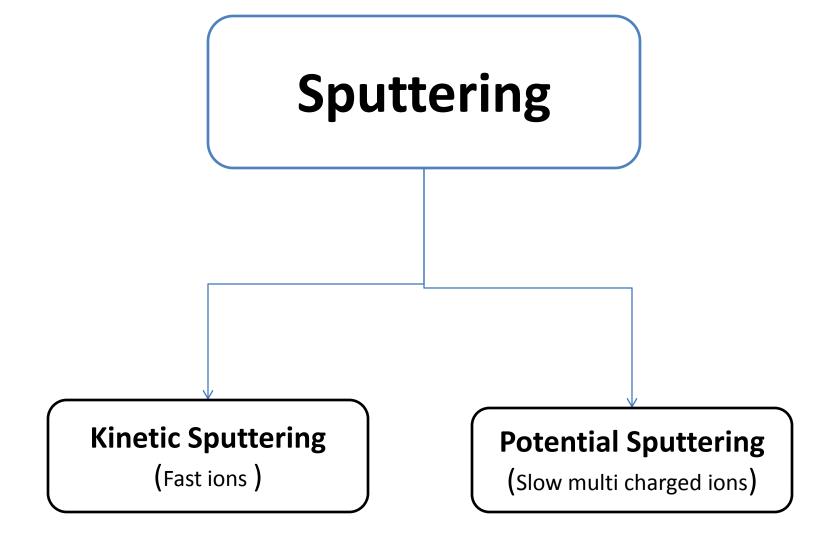
### Introduction

 At energies around 1 keV/amu, SW protons and HI interact with the lunar surface materials via a number of microscopic interactions, but for our purposes, the most important of these is atomic sputtering

# **Physics of Sputtering**

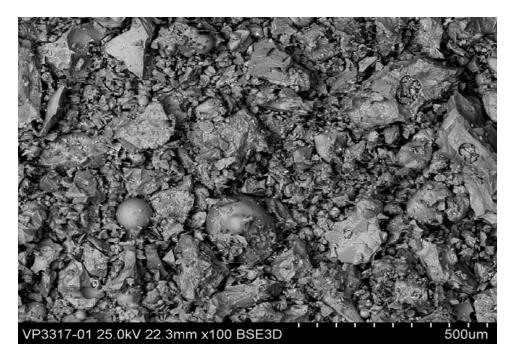
When the cascade gives the target atom energy greater than the surface binding energy, then the atom may be sputtered





#### Lunar Regolith Simulant JSC-1A AGGL

XPS: Surface of the simulant consists mostly of oxides



Element	C	0	Si	Al	Fe	Ca	Mg	Ti	Na	Р	K	Cr	F
Atomic %	2.3	55.6	19.5	8.4	1.4	4.3	3.9	0.4	3.3	0.3	0.3	0.1	0.1

#### **Heat-Conduction/Diffusion Equation**

$$\frac{1}{\kappa}\frac{\partial u}{\partial t} = \nabla^2 u$$

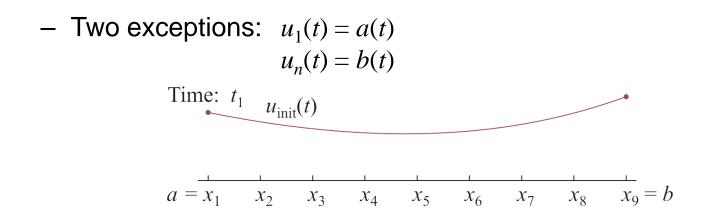
Discretizing the space and time components in one dimension gives:

$$\frac{u(x,t+h)-u(x,t)}{\Delta t} = \frac{\kappa}{h^2} \left( u(x-h,t) - 2u(x,t) + u(x+h,t) \right)$$
  
which allowed us to find

$$u(x,t+h) = u(x,t) + \frac{\kappa \Delta t}{h^2} \left( u(x-h,t) - 2u(x,t) + u(x+h,t) \right)$$

# Method of Lines

As an alternative approach, associate with each spatial point an unknown function  $u_k(t)$ 



This approach was popularized by the chemical engineer William E. Schiesser in his 1991 text *The Numerical Method* of Lines

# Method of Lines

In order to substitute  $u_k(t)$  into our mixed partial-/finite-difference equation,

$$\frac{\partial}{\partial t}u(x,t) = \frac{\kappa}{h^2} \left( u(x-h,t) - 2u(x,t) + u(x+h,t) \right)$$

we note that the solution at location x - h is  $u_{k-1}(t)$ and the solution at x + h is  $u_{k+1}(t)$ :

$$\frac{d}{dt}u_{k}(t) = \frac{\kappa}{h^{2}}\left(u_{k-1}(t) - 2u_{k}(t) + u_{k+1}(t)\right)$$

We also have the initial condition:

 $u_k(t_{\text{initial}}) = u_{\text{init}}(x_k)$ 

## Systems of IVPs

We can therefore write this as:

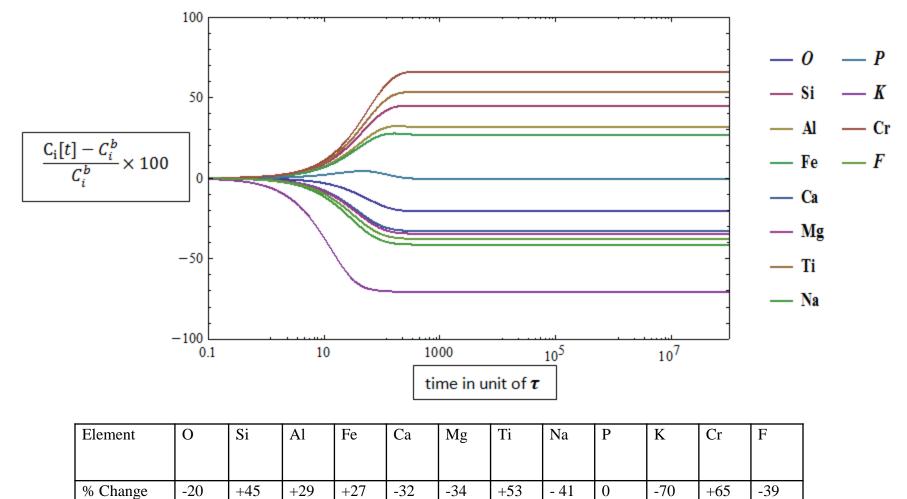
$$\mathbf{u}^{(1)}(t) = \mathbf{f}(t, \mathbf{u}(t)) = \frac{\kappa}{h^2} \begin{pmatrix} -2 & 1 & & \\ 1 & -2 & 1 & \\ & 1 & -2 & \ddots & \\ & & \ddots & \ddots & 1 \\ & & & 1 & -2 \end{pmatrix} \mathbf{u}(t) + \begin{pmatrix} a(t) \\ 0 \\ \vdots \\ 0 \\ b(t) \end{pmatrix} \end{pmatrix}$$

where 
$$\mathbf{f}(t, \mathbf{u}) \stackrel{def}{=} \frac{\kappa}{h^2} \left( \begin{pmatrix} -2 & 1 & & \\ 1 & -2 & 1 & \\ & 1 & -2 & \ddots & \\ & & \ddots & \ddots & 1 \\ & & & \ddots & \ddots & 1 \\ & & & & 1 & -2 \end{pmatrix} \mathbf{u} + \begin{pmatrix} a(t) \\ 0 \\ 0 \\ 0 \\ b(t) \end{pmatrix} \right)$$

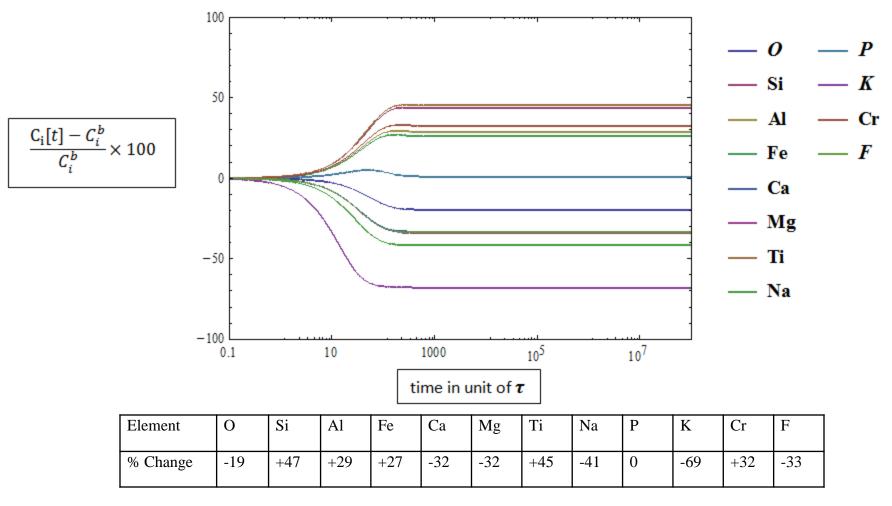
# **Non-Equilibrium Model** $\frac{dC_i}{dt} = \frac{1}{\tau} \left[ -C_i \sum_j Y_{ij} f_j + C_i^b \sum_k C_k Y_{kj} f_j \right]$

- C<sub>i</sub> is the abundant of element i in JSC-1A AGGL
- C<sup>b</sup><sub>i</sub> is the fractional abundant of element i in the JSC-1A AGGL bulk
- $Y_{ij}$  is the yield of element i by solar wind ion j,
- F<sub>j</sub> is the fraction of solar wind j in the solar wind flux
- T is a constant has dimension of time.

Calculated changes in the elemental composition of a JSC-1A AGGL surface as a function of time due to the kinetic sputtering of the solar-wind protons.



Calculated changes in the elemental composition of a JSC-1A AGGL surface as a function of time due to the kinetic sputtering of the solar-wind protons and heavy ions.



## Conclusions

- Sputtering is an important dynamic mechanism that affects the composition of both the lunar surface and its tenuous exosphere.
- The contribution of the solar-wind protons and ions kinetic sputtering to the changes in the composition of the surface layer of the oxides of the lunar surface is well understood and modeled
- we expect these changes to be more than the calculated due to contribution of the potential sputtering which is unclear.

### Conclusions

 The changes in the elemental abundant of JSC-1A AGGL due to the solar wind protons alone approach the steady state for times close to 300T, and this time shortened to about 200 T due to the heavy ions contributions